## Transport Phenomena in Food Processes Academic Year 2012-2013

| Last name | First name | ID | Signature |
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Problem 1. A power-law fluid (with constant fluid properties, $\rho, m, n$ ) flows through a tapered tube with circular cross-section, for which the tube radius changes linearly from $R_{0}$ at the tube entrance to a slightly smaller value $R_{L}$ at the tube exit.

1. Applying locally the result obtained for a circular tube (eq. 8.3-9), write the proper differential equation in $\mathrm{d} \mathcal{P} / \mathrm{d} z$;
2. Integrate the equation obtained in 1 . in order to get the pressure gradient;
3. Compare the mass flow rate obtainable for the power-law fluid and for an incompressible fluid of constant properties $(\rho, \mu)$, flowing in the same tube in presence of the same pressure gradient (the solution for incompressible Newtonian fluid is given by eq. 2B.10-3).

Data. $L=10 \mathrm{~m}, \rho=1000 \mathrm{~kg} / \mathrm{m}^{3}, m=0.2 \mathrm{~Pa} \cdot \mathrm{~s}^{n}, n=0.85, R_{0}=10 \mathrm{~cm}, R_{L}=9 \mathrm{~cm}, \Delta \mathcal{P} / L=1.4 \mathrm{bar} / \mathrm{m}$, $\mu=0.2 \mathrm{~Pa} \cdot \mathrm{~s}$.

Problem 2. A piece of foodstuff shaped as a long cylinder with radius $R$ and constant properties ( $k, \rho \hat{C}_{P}$ ) was taken from a fridge with an initial uniform temperature $T_{0}$, and it was heated by hot air at temperature $T_{\infty}$, flowing orthogonally to the cylinder axis, with velocity $v_{\infty}$. Under these conditions, the Churchill \& Bernstein correlation holds, and the interphase heat transfer coefficient has the value $h$. After a time $t^{*}$, the axial temperature of the foodstuff was measured to be $T_{A}$. Calculate:

1. the value of air velocity, $v_{\infty}$;
2. the foodstuff thermal conductivity, $k$;
3. the surface temperature after the time $t^{*}, T_{s}\left(t^{*}, r=R\right)$.

The Churchill and Bernstein correlation is ( $D$ being the cylinder diameter):

$$
N_{N u}=\frac{h D}{k}=0.3+\frac{0.62 N_{R e}^{0.5} N_{P r}^{0.33}}{\left[1+\left(\frac{0.4}{N_{P r}}\right)^{0.67}\right]^{0.25}}\left[1+\left(\frac{N_{R e}}{282000}\right)^{0.625}\right]^{0.8}
$$

The air properties could be taken as constants on their initial values.
Data. $R=5 \mathrm{~cm}, \rho \hat{C}_{P}=6.0 \mathrm{MJ} /\left(\mathrm{m}^{3} \mathrm{~K}\right), T_{0}=5^{\circ} \mathrm{C}, T_{\infty}=150^{\circ} \mathrm{C}, h=35 \mathrm{~W} /\left(\mathrm{m}^{2} \mathrm{~K}\right), t^{*}=0.5 \mathrm{hr}, T_{A}=30^{\circ} \mathrm{C}$.

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$$
\mathrm{L}_{\mathrm{m}}:=10 \cdot \mathrm{~m} \quad \rho:=1000 \frac{\mathrm{~kg}}{\mathrm{~m}^{3}} \quad \mathrm{~m}:=0.2 \cdot \mathrm{~Pa} \cdot \mathrm{~s} .05 \quad \mathrm{n}:=0.85 \quad \mathrm{R}_{0}:=10 \cdot \mathrm{~cm} \quad \mathrm{R}_{\mathrm{L}}:=9 \cdot \mathrm{~cm} \quad \Delta \mathrm{P}_{\mathrm{L}} \mathrm{~L}:=1.4 \frac{\mathrm{~Pa}}{\mathrm{~m}} \quad \mu:=0.2 \cdot \mathrm{~Pa} \cdot \mathrm{~s}
$$

We consider a small region of the tapered tube to be a straight tube over a short distance $d z$; then we can write "locally"

$$
\begin{equation*}
w=\frac{\pi R^{3} \rho}{(1 / n)+3}\left[-\frac{d \mathcal{P}}{d z} \frac{R}{2 m}\right]^{1 / n} \tag{8.3-9}
\end{equation*}
$$

$$
d R \_d z:=\frac{R_{L}-R_{0}}{L}=-1 \times 10^{-3}
$$

Take the $n$th power of both sides to get

$$
-\frac{d \rho^{\rho}}{d z}=\frac{2 m}{R}\left[\frac{w}{\pi R^{3} \rho}\left(\frac{1}{n}+3\right)\right]^{n}
$$

$$
\frac{\mathrm{d}}{\mathrm{dz}} \mathrm{P}=-\frac{2 \cdot \mathrm{~m}}{\mathrm{R}} \cdot\left[\frac{\mathrm{w}}{\pi \cdot \mathrm{R}^{3} \cdot \rho} \cdot\left(\frac{1}{\mathrm{n}}+3\right)\right]^{\mathrm{n}}
$$

in which $R$ is a function of $z$ :

$$
R=R_{0}+\left(\frac{R_{L}-R_{0}}{L}\right) z
$$

It is easier to integrate the differential equation if we rewrite it as

$$
-\frac{d \mathscr{P}}{d R} \frac{d R}{d z}=-\frac{d \mathcal{P}}{d R}\left(\frac{R_{L}-R_{0}}{L}\right)=\frac{2 m}{R}\left[\frac{w}{\pi R^{3} \rho}\left(\frac{1}{n}+3\right)\right]^{n}
$$

Then when this equation is integrated with respect to $R$, we get

$$
-\int_{\mathfrak{P}_{0}}^{\mathfrak{P}_{L}} d \mathscr{P}=\left(\frac{2 m L}{R_{L}-R_{0}}\right)\left[\frac{w}{\pi \rho}\left(\frac{1}{n}+3\right)\right]^{n} \int_{R_{0}}^{R_{L}} \frac{1}{R^{3 n+1}} d R
$$

## Therefore

$$
\begin{aligned}
\wp_{0}-\wp_{L} & =\left(\frac{2 m L}{R_{L}-R_{0}}\right)\left[\frac{w}{\pi \rho}\left(\frac{1}{n}+3\right)\right]^{n}\left(\frac{R_{L}^{-3 n}-R_{0}^{-3 n}}{-3 n}\right) \\
& =\left(\frac{2 m L}{3 n}\right)\left[\frac{w}{\pi \rho}\left(\frac{1}{n}+3\right)\right]^{n}\left(\frac{R_{L}^{-3 n}-R_{0}^{-3 n}}{R_{0}-R_{L}}\right)
\end{aligned}
$$

$$
\frac{\Delta \mathrm{P}}{\mathrm{~L}}=\frac{2 \cdot \mathrm{~m}}{3 \cdot \mathrm{n}} \cdot\left[\frac{\mathrm{w}}{\pi \cdot \rho} \cdot\left(\frac{1}{\mathrm{n}}+3\right)\right]^{\mathrm{n}} \cdot\left(\frac{\mathrm{R}_{\mathrm{L}}^{-3 \cdot \mathrm{n}}-\mathrm{R}_{0}^{-3 \cdot \mathrm{n}}}{\mathrm{R}_{0}-\mathrm{R}_{\mathrm{L}}}\right)
$$

$\mathrm{w}:=\frac{\pi \cdot \rho}{\frac{1}{\mathrm{n}}+3} \cdot\left[\frac{3 \cdot \mathrm{n}}{2 \cdot \mathrm{~m}} \cdot \Delta \mathrm{P}_{-} \mathrm{L} \cdot\left(\frac{\mathrm{R}_{0}-\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{\mathrm{L}}^{-3 \cdot 0.85}-\mathrm{R}_{0}-3 \cdot 0.85}\right)\right]^{\frac{1}{0.85}}=0.175 \frac{\mathrm{~kg}}{\mathrm{~s}}$
$w=\frac{\pi\left(\mathscr{P}_{0}-\mathscr{P}_{L}\right) R_{0}^{4} \rho}{8 \mu L}\left[1-\frac{1+\left(R_{L} / R_{0}\right)+\left(R_{L} / R_{0}\right)^{2}-3\left(R_{L} / R_{0}\right)^{3}}{1+\left(R_{L} / R_{0}\right)+\left(R_{L} / R_{0}\right)^{2}}\right]$
$\mathrm{w}_{\mu}:=\frac{\pi \cdot \rho \cdot \Delta \mathrm{P}_{-} \mathrm{L} \cdot \mathrm{R}_{0}}{8 \cdot \mu} \cdot\left[1-\frac{1+\frac{\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{0}}+\left(\frac{\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{0}}\right)^{2}-3 \cdot\left(\frac{\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{0}}\right)^{3}}{1+\frac{\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{0}}+\left(\frac{\mathrm{R}_{\mathrm{L}}}{\mathrm{R}_{0}}\right)^{2}}\right]=0.222 \frac{\mathrm{~kg}}{\mathrm{~s}}$

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$$
\begin{aligned}
& \mathrm{T}_{\text {inf }}:=150^{\circ} \mathrm{C} \quad \mathrm{~T}_{0}:=5^{\circ} \mathrm{C} \quad \mathrm{~T}_{\mathrm{A}}:=30^{\circ} \mathrm{C} \quad \mathrm{R}:=5 \cdot \mathrm{~cm} \quad \mathrm{~h}:=35 \cdot \frac{\mathrm{~W}}{\mathrm{~m}^{2} \cdot \mathrm{~K}} \quad \mathrm{t}:=\frac{1}{2} \cdot \mathrm{hr} \\
& \rho \mathrm{C}_{\mathrm{P}}:=3000 \cdot \frac{\mathrm{~kg}}{\mathrm{~m}^{3}} \cdot 2000 \cdot \frac{\mathrm{~J}}{\mathrm{~kg} \cdot \mathrm{~K}}=6 \times 10^{6} \cdot \frac{\mathrm{~J}}{\mathrm{~m}^{3} \cdot \mathrm{~K}} \\
& \mathrm{~N}_{\mathrm{Nu}}\left(\mathrm{~N}_{\mathrm{Re}}, \mathrm{~N}_{\mathrm{Pr}}\right):=0.3+\frac{0.62 \cdot \mathrm{~N}_{\mathrm{Re}}{ }^{0.5} \cdot \mathrm{~N}_{\mathrm{Pr}}}{\left[1+\left(\frac{0.4}{\mathrm{~N}_{\mathrm{Pr}}}\right)^{0.33}\right.}\left[1+\left(\frac{\mathrm{N}_{\mathrm{Re}}}{282000}\right)^{0.625}\right]^{0.8}
\end{aligned}
$$

1. The Churchill \& Bernstein correlation has to be solved iteratively in the unknown v.inf

$$
\begin{aligned}
& \mathrm{T}_{\mathrm{f}}:=\frac{\mathrm{T}_{\mathrm{inf}}+\mathrm{T}_{0}}{2}=77.5 \cdot{ }^{\circ} \mathrm{C} \quad \mathrm{~N}_{\operatorname{Pr} \cdot \mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)=0.706 \quad \mathrm{k}_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)=0.03 \cdot \frac{\mathrm{~W}}{\mathrm{~m} \cdot \mathrm{~K}} \quad v_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)=2.068 \times 10^{-5} \frac{5 \mathrm{~m}^{2}}{\mathrm{~s}} \\
& \mathrm{v}_{\mathrm{inf}}:=1 \cdot \frac{\mathrm{~m}}{\mathrm{~s}} \quad \text { Given } \quad \frac{\mathrm{h} \cdot 2 \cdot \mathrm{R}}{\mathrm{k}_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)}=\mathrm{N}_{\mathrm{Nu}}\left(\frac{\mathrm{v}_{\mathrm{inf}} \cdot 2 \cdot \mathrm{R}}{\nu_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)}, \mathrm{N}_{\mathrm{Pr} . \mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)\right) \quad \mathrm{v}_{\mathrm{inf}}:=\operatorname{Minerr}\left(\mathrm{v}_{\mathrm{inf}}\right)=8.052 \frac{\mathrm{~m}}{\mathrm{~s}} \\
& \mathrm{~N}_{\mathrm{Re}}:=\frac{\mathrm{v}_{\mathrm{inf}} \cdot 2 \cdot \mathrm{R}}{\nu_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)}=3.893 \times 10^{4} \quad \quad \mathrm{~N}_{\mathrm{Nu}}\left(\mathrm{~N}_{\mathrm{Re}}, \mathrm{~N}_{\mathrm{Pr} . \mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)\right)=117.67 \\
& \mathrm{j}_{\mathrm{H}}\left(\frac{\mathrm{v}_{\mathrm{inf}} \cdot 2 \cdot \mathrm{R}}{\nu_{\mathrm{A}}\left(\mathrm{~T}_{\mathrm{f}}\right)}, \mathrm{N}_{\text {Pr.A }}\left(\mathrm{T}_{\mathrm{f}}\right)\right)=3.391 \times 10^{-3} \\
& \text { 2. The transient heating of a cylinder can be solved } \\
& \text { under the single-term approximation (to be checked } \\
& \text { later) }
\end{aligned}
$$ later)

$$
\begin{equation*}
\lambda_{1} \cdot \frac{\mathrm{~J}_{1}\left(\lambda_{1}\right)}{\mathrm{J}_{0}\left(\lambda_{1}\right)}=\frac{\mathrm{h} \cdot \mathrm{R}}{\mathrm{k}} \tag{1}
\end{equation*}
$$

(2) $\quad A_{1}=\frac{2}{\lambda_{1}} \cdot \frac{J_{1}\left(\lambda_{1}\right)}{J_{1}\left(\lambda_{1}\right)^{2}+J_{0}\left(\lambda_{1}\right)^{2}}$

$$
\begin{equation*}
\theta_{A}=\theta(t, r=0)=A_{1} \cdot J_{0}\left(\lambda_{1} \cdot 0\right) \cdot \exp \left(-\lambda_{1}{ }^{2} \cdot \tau\right)=A_{1} \cdot J_{0}\left(\lambda_{1} \cdot 0\right) \cdot \exp \left(-\lambda_{1}{ }^{2} \cdot \frac{\mathrm{k}}{\rho \mathrm{C}_{\mathrm{P}}} \cdot \frac{\mathrm{t}}{\mathrm{R}^{2}}\right) \tag{3}
\end{equation*}
$$

from (1)

$$
\begin{equation*}
\mathrm{k}=\frac{\mathrm{h} \cdot \mathrm{R} \cdot \mathrm{~J}_{0}\left(\lambda_{1}\right)}{\lambda_{1} \cdot \mathrm{~J}_{1}\left(\lambda_{1}\right)} \tag{1'}
\end{equation*}
$$

substituting (1') and (2) in (3), one equation (3') in one unknown ( $\lambda .1$ ) is obtained

$$
\begin{aligned}
& \theta_{\mathrm{A}}=\frac{2}{\lambda_{1}} \cdot \frac{\mathrm{~J}_{1}\left(\lambda_{1}\right)}{\mathrm{J}_{1}\left(\lambda_{1}\right)^{2}+\mathrm{J}_{0}\left(\lambda_{1}\right)^{2}} \cdot \mathrm{~J}_{0}\left(\lambda_{1} \cdot 0\right) \cdot \exp \left(-\lambda_{1}{ }^{2} \cdot \frac{\mathrm{~h} \cdot \mathrm{R} \cdot \mathrm{~J}_{0}\left(\lambda_{1}\right)}{\lambda_{1} \cdot \mathrm{~J}_{1}\left(\lambda_{1}\right)} \cdot \frac{1}{\rho \mathrm{C}_{\mathrm{P}}} \cdot \frac{\mathrm{t}}{\mathrm{R}^{2}}\right) \quad\left(3^{\prime}\right) \\
& \lambda_{1}:=1 \quad \text { Given } \quad \theta_{\mathrm{A}}=\frac{2}{\lambda_{1}} \cdot \frac{\mathrm{~J} 1\left(\lambda_{1}\right)}{\mathrm{J} 1\left(\lambda_{1}\right)^{2}+\mathrm{J} 0\left(\lambda_{1}\right)^{2}} \cdot \mathrm{~J} 0\left(\lambda_{1} \cdot 0\right) \cdot \exp \left(-\lambda_{1}{ }^{2} \cdot \frac{\mathrm{~h} \cdot \mathrm{R} \cdot \mathrm{~J} 0\left(\lambda_{1}\right)}{\lambda_{1} \cdot \mathrm{~J} 1\left(\lambda_{1}\right)} \cdot \frac{1}{\rho \mathrm{C}_{\mathrm{P}}} \cdot \frac{\mathrm{t}}{\mathrm{R}^{2}}\right) \\
& \qquad \lambda_{1}:=\operatorname{Minerr}\left(\lambda_{1}\right)=1.144 \\
& \text { Then, the thermal conductivity could be calculated by eq. (1') } \quad \mathrm{k}:=\frac{\mathrm{h} \cdot \mathrm{R} \cdot \mathrm{~J} 0\left(\lambda_{1}\right)}{\lambda_{1} \cdot \mathrm{JJ}\left(\lambda_{1}\right)}=2.209 \cdot \frac{\mathrm{~W}}{\mathrm{~m} \cdot \mathrm{~K}}
\end{aligned}
$$

Biot's number $\quad \mathrm{N}_{\mathrm{Bi}}:=\frac{\mathrm{h} \cdot \mathrm{R}}{\mathrm{k}}=0.792$
Fourier's number $\quad \tau:=\frac{\mathrm{k}}{\rho \mathrm{C}_{\mathrm{P}}} \cdot \frac{\mathrm{t}}{\mathrm{R}^{2}}=0.265 \quad$ being higher than 0.2 , the single-term approximation holds.
3. To calculate the surface temperature, apply the single-term approximation for a long cylinder at the surface ( $\xi=1$ )

$$
\begin{aligned}
& \theta_{\mathrm{S}}:=\frac{2}{\lambda_{1}} \cdot \frac{\mathrm{~J} 1\left(\lambda_{1}\right)}{\mathrm{J} 1\left(\lambda_{1}\right)^{2}+\mathrm{J} 0\left(\lambda_{1}\right)^{2}} \cdot \mathrm{~J} 0\left(\lambda_{1} \cdot 1\right) \cdot \exp \left(-\lambda_{1}{ }^{2} \cdot \tau\right)=0.578 \\
& \mathrm{~T}_{\mathrm{S}}:=\mathrm{T}_{\mathrm{inf}}+\left(\mathrm{T}_{0}-\mathrm{T}_{\mathrm{inf}}\right) \cdot \theta_{\mathrm{S}}=66.193 \cdot{ }^{\circ} \mathrm{C}
\end{aligned}
$$

