Last name	First name	ID	Signature

Problem 1. A power-law fluid (with constant fluid properties, ρ , m, n) flows through a tapered tube with circular cross-section, for which the tube radius changes linearly from R_0 at the tube entrance to a slightly smaller value R_L at the tube exit.

- 1. Applying locally the result obtained for a circular tube (eq. 8.3-9), write the proper differential equation in $d\mathcal{P}/dz$;
- 2. Integrate the equation obtained in 1. in order to get the pressure gradient;
- 3. Compare the mass flow rate obtainable for the power-law fluid and for an incompressible fluid of constant properties (ρ , μ), flowing in the same tube in presence of the same pressure gradient (the solution for incompressible Newtonian fluid is given by eq. 2B.10-3).

Data. L = 10 m, $\rho = 1000 \text{ kg/m}^3$, $m = 0.2 \text{ Pa} \cdot \text{s}^n$, n = 0.85, $R_0 = 10 \text{ cm}$, $R_L = 9 \text{ cm}$, $\Delta \mathcal{P}/L = 1.4 \text{ bar/m}$, $\mu = 0.2 \text{ Pa} \cdot \text{s}$.

Problem 2. A piece of foodstuff shaped as a long cylinder with radius *R* and constant properties $(k, \rho \hat{C}_P)$ was taken from a fridge with an initial uniform temperature T_0 , and it was heated by hot air at temperature T_{∞} , flowing orthogonally to the cylinder axis, with velocity v_{∞} . Under these conditions, the Churchill & Bernstein correlation holds, and the interphase heat transfer coefficient has the value *h*. After a time t^* , the axial temperature of the foodstuff was measured to be T_A . Calculate:

- 1. the value of air velocity, v_{∞} ;
- 2. the foodstuff thermal conductivity, *k*;
- 3. the surface temperature after the time t^* , $T_s(t^*, r = R)$.

The Churchill and Bernstein correlation is (*D* being the cylinder diameter):

$$N_{Nu} = \frac{hD}{k} = 0.3 + \frac{0.62N_{Re}^{0.5}N_{Pr}^{0.33}}{\left[1 + \left(\frac{0.4}{N_{Pr}}\right)^{0.67}\right]^{0.25}} \left[1 + \left(\frac{N_{Re}}{282000}\right)^{0.625}\right]^{0.8}$$

The air properties could be taken as constants on their initial values.

Data. R = 5 cm, $\rho \hat{C}_P = 6.0 \text{ MJ/(m^3K)}$, $T_0 = 5^{\circ}\text{C}$, $T_{\infty} = 150^{\circ}\text{C}$, $h = 35 \text{ W/(m^2K)}$, $t^* = 0.5 \text{ hr}$, $T_A = 30^{\circ}\text{C}$.

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 $\underset{\mathsf{m}}{\text{L}} \coloneqq 10 \text{ m} \quad \rho \coloneqq 1000 \frac{\text{kg}}{\text{m}^3} \quad \mathsf{m} \coloneqq 0.2 \cdot \text{Pa} \cdot \text{s}^{0.85} \quad \mathsf{n} \coloneqq 0.85 \qquad \mathsf{R}_0 \coloneqq 10 \text{ cm} \quad \mathsf{R}_L \coloneqq 9 \cdot \text{cm} \quad \Delta P_L \coloneqq 1.4 \cdot \frac{\text{Pa}}{\text{m}} \quad \mu \coloneqq 0.2 \cdot \text{Pa} \cdot \text{s}^{0.85}$

We consider a small region of the tapered tube to be a straight tube over a short distance dz; then we can write "locally"

$$w = \frac{\pi R^{3} \rho}{(1/n) + 3} \left[-\frac{d\mathbf{P}}{dz} \frac{R}{2m} \right]^{1/n} \qquad \text{eq. (8.3-9)}$$

Take the *n*th power of both sides to get

$$-\frac{d\mathbf{\mathcal{P}}}{dz} = \frac{2m}{R} \left[\frac{w}{\pi R^3 \rho} \left(\frac{1}{n} + 3 \right) \right]^n$$

in which *R* is a function of *z*:

 $R = R_0 + \left(\frac{R_L - R_0}{L}\right)z$

It is easier to integrate the differential equation if we rewrite it as

$$-\frac{d\mathbf{\mathcal{P}}}{dR}\frac{dR}{dz} = -\frac{d\mathbf{\mathcal{P}}}{dR}\left(\frac{R_L - R_0}{L}\right) = \frac{2m}{R}\left[\frac{w}{\pi R^3 \rho}\left(\frac{1}{n} + 3\right)\right]^n$$

Then when this equation is integrated with respect to *R*, we get

$$-\int_{\mathcal{P}_0}^{\mathcal{P}_L} d\mathcal{P} = \left(\frac{2mL}{R_L - R_0}\right) \left[\frac{w}{\pi\rho} \left(\frac{1}{n} + 3\right)\right]^n \int_{R_0}^{R_L} \frac{1}{R^{3n+1}} dR$$

Therefore

$$\mathbf{\mathcal{P}}_{0} - \mathbf{\mathcal{P}}_{L} = \left(\frac{2mL}{R_{L} - R_{0}}\right) \left[\frac{w}{\pi\rho} \left(\frac{1}{n} + 3\right)\right]^{n} \left(\frac{R_{L}^{-3n} - R_{0}^{-3n}}{-3n}\right)$$
$$= \left(\frac{2mL}{3n}\right) \left[\frac{w}{\pi\rho} \left(\frac{1}{n} + 3\right)\right]^{n} \left(\frac{R_{L}^{-3n} - R_{0}^{-3n}}{R_{0} - R_{L}}\right)$$

	(-3.n)	3. n
ΔΡ	$2 \cdot m [w] (1) = \sqrt{n} R_{I} = R_{0}$	5 11
<u> </u>	$\frac{2}{10} \cdot \frac{1}{10} \cdot \frac{1}{10} + 3 \cdot \frac{1}{10} \cdot \frac{1}{100} \cdot \frac{1}$	
L	$3 \cdot n = \pi \cdot 0 \cdot n$	
-		

d _P	2·m	w	(1	$\begin{bmatrix} 2 \end{bmatrix}^n$
dz	R	$\pi \cdot R^3 \cdot \rho$	(<u>n</u>	

 $dR_dz := \frac{R_L - R_0}{L} = -1 \times 10^{-3}$

$$\mathbf{w} := \frac{\pi \cdot \rho}{\frac{1}{n} + 3} \cdot \left[\frac{3 \cdot n}{2 \cdot m} \cdot \Delta \mathbf{P}_{\mathbf{L}} \cdot \left(\frac{\mathbf{R}_{0} - \mathbf{R}_{\mathbf{L}}}{\mathbf{R}_{\mathbf{L}}^{-3 \cdot 0.85} - \mathbf{R}_{0}^{-3 \cdot 0.85}} \right) \right]^{\frac{1}{0.85}} = 0.175 \frac{\text{kg}}{\text{s}}$$

$$w_{\mu} \coloneqq \frac{\pi(\mathcal{P}_{0} - \mathcal{P}_{L})R_{0}^{4}\rho}{8\mu L} \left[1 - \frac{1 + (R_{L}/R_{0}) + (R_{L}/R_{0})^{2} - 3(R_{L}/R_{0})^{3}}{1 + (R_{L}/R_{0}) + (R_{L}/R_{0})^{2}} \right]$$
(2B.10-3)
$$w_{\mu} \coloneqq \frac{\pi \cdot \rho \cdot \Delta P_{\perp} L \cdot R_{0}^{4}}{8 \cdot \mu} \cdot \left[1 - \frac{1 + \frac{R_{L}}{R_{0}} + \left(\frac{R_{L}}{R_{0}}\right)^{2} - 3 \cdot \left(\frac{R_{L}}{R_{0}}\right)^{3}}{1 + \frac{R_{L}}{R_{0}} + \left(\frac{R_{L}}{R_{0}}\right)^{2}} \right] = 0.222 \frac{kg}{s}$$

Problem 2. A piece of foodstuff shaped as a long cylinder with radius *R* and constant properties $(k, \rho \hat{C}_P)$ was taken from a fridge with an initial uniform temperature T_0 , and it was heated by hot air at temperature T_∞ , flowing orthogonally to the cylinder axis, with velocity v_∞ . Under these conditions, the Churchill & Bernstein correlation holds, and the interphase heat transfer coefficient has the value *h*. After a time t^* , the axial temperature of the foodstuff was measured to be T_A . Calculate:

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- 2. the foodstuff thermal conductivity, k;
- the surface temperature after the time t*, T_s(t*, r = R).

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The air properties could be taken as constants on their initial values.

Data. $R = 5 \text{ cm}, \rho \hat{C}_P = 6.0 \text{ MJ}/(\text{m}^3\text{K}), T_0 = 5^{\circ}\text{C}, T_{\infty} = 150^{\circ}\text{C}, h = 35 \text{ W}/(\text{m}^2\text{K}), t^* = 0.5 \text{ hr}, T_A = 30^{\circ}\text{C}.$

$$T_{inf} := 150 \,^{\circ}C \qquad T_0 := 5 \,^{\circ}C \qquad T_A := 30 \,^{\circ}C \qquad R := 5 \,^{\circ}cm \qquad h := 35 \,^{\circ}\frac{w}{m^2 \cdot K} \qquad t := \frac{1}{2} \cdot hr$$

$$\rho C_P := 3000 \cdot \frac{kg}{m^3} \cdot 2000 \cdot \frac{J}{kg \cdot K} = 6 \times 10^6 \cdot \frac{J}{m^3 \cdot K}$$

$$N_{Nu} (N_{Re}, N_{Pr}) := 0.3 + \frac{0.62 \cdot N_{Re}}{\left[1 + \left(\frac{0.4}{N_{Pr}}\right)^{0.67}\right]^{0.25}} \cdot \left[1 + \left(\frac{N_{Re}}{282000}\right)^{0.625}\right]^{0.8}$$

1. The Churchill & Bernstein correlation has to be solved iteratively in the unknown v.inf

$$T_{f} \coloneqq \frac{T_{inf} + T_{0}}{2} = 77.5 \cdot \circ C \qquad N_{Pr,A}(T_{f}) = 0.706 \qquad k_{A}(T_{f}) = 0.03 \cdot \frac{W}{m \cdot K} \qquad \nu_{A}(T_{f}) = 2.068 \times 10^{-5} \frac{m^{2}}{s}$$

$$v_{inf} \coloneqq 1 \cdot \frac{m}{s} \qquad \text{Given} \qquad \frac{h \cdot 2 \cdot R}{k_{A}(T_{f})} = N_{Nu} \left(\frac{v_{inf} \cdot 2 \cdot R}{\nu_{A}(T_{f})}, N_{Pr,A}(T_{f}) \right) \qquad v_{inf} \coloneqq \text{Minerr}(v_{inf}) = 8.052 \frac{m}{s}$$

$$N_{Re} \coloneqq \frac{v_{inf} \cdot 2 \cdot R}{\nu_{A}(T_{f})} = 3.893 \times 10^{4} \qquad N_{Nu}(N_{Re}, N_{Pr,A}(T_{f})) = 117.67$$

$$j_{H} \left(\frac{v_{inf} \cdot 2 \cdot R}{\nu_{A}(T_{f})}, N_{Pr,A}(T_{f}) \right) = 3.391 \times 10^{-3}$$

2. The transient heating of a cylinder can be solved under the single-term approximation (to be checked later)

$$\theta_{\rm A} := \frac{{\rm T_{inf}} - {\rm T_{\rm A}}}{{\rm T_{inf}} - {\rm T_{\rm 0}}} = 0.828$$

(1)
$$\lambda_1 \cdot \frac{J_1(\lambda_1)}{J_0(\lambda_1)} = \frac{h \cdot R}{k}$$

(2)
$$A_1 = \frac{2}{\lambda_1} \cdot \frac{J_1(\lambda_1)}{J_1(\lambda_1)^2 + J_0(\lambda_1)^2}$$

(3)
$$\theta_{A} = \theta(t, r = 0) = A_{1} \cdot J_{0}(\lambda_{1} \cdot 0) \cdot \exp(-\lambda_{1}^{2} \cdot \tau) = A_{1} \cdot J_{0}(\lambda_{1} \cdot 0) \cdot \exp\left(-\lambda_{1}^{2} \cdot \frac{k}{\rho C_{p}} \cdot \frac{t}{R^{2}}\right)$$

from (1)
$$k = \frac{h \cdot R \cdot J_{0}(\lambda_{1})}{\lambda_{1} \cdot J_{1}(\lambda_{1})}$$
(1')

substituting (1') and (2) in (3), one equation (3') in one unknown (λ .1) is obtained

$$\theta_{A} = \frac{2}{\lambda_{1}} \cdot \frac{J_{1}(\lambda_{1})}{J_{1}(\lambda_{1})^{2} + J_{0}(\lambda_{1})^{2}} \cdot J_{0}(\lambda_{1} \cdot 0) \cdot \exp\left(-\lambda_{1}^{2} \cdot \frac{h \cdot R \cdot J_{0}(\lambda_{1})}{\lambda_{1} \cdot J_{1}(\lambda_{1})} \cdot \frac{1}{\rho C_{p}} \cdot \frac{t}{R^{2}}\right) \quad (3')$$

$$\lambda_{1} := 1 \qquad \text{Given} \qquad \theta_{A} = \frac{2}{\lambda_{1}} \cdot \frac{J_{1}(\lambda_{1})}{J_{1}(\lambda_{1})^{2} + J_{0}(\lambda_{1})^{2}} \cdot J_{0}(\lambda_{1} \cdot 0) \cdot \exp\left(-\lambda_{1}^{2} \cdot \frac{h \cdot R \cdot J_{0}(\lambda_{1})}{\lambda_{1} \cdot J_{1}(\lambda_{1})} \cdot \frac{1}{\rho C_{p}} \cdot \frac{t}{R^{2}}\right)$$

$$\lambda_{1} := Minor(\lambda_{1}) = 1.144$$

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Then, the thermal conductivity could be calculated by eq. (1')

Biot's number

$$N_{Bi} := \frac{h \cdot R}{k} = 0.792$$

 $\tau := \frac{k}{\rho C_P} \cdot \frac{t}{R^2} = 0.265$

Fourier's number

$$\mathbf{k} := \frac{\mathbf{h} \cdot \mathbf{R} \cdot \mathbf{J0}(\lambda_1)}{\lambda_1 \cdot \mathbf{J1}(\lambda_1)} = 2.209 \cdot \frac{\mathbf{W}}{\mathbf{m} \cdot \mathbf{K}}$$

being higher than 0.2, the single-term approximation holds.

3. To calculate the surface temperature, apply the single-term approximation for a long cylinder at the surface ($\xi = 1$)

$$\theta_{\mathbf{S}} \coloneqq \frac{2}{\lambda_{1}} \cdot \frac{J1(\lambda_{1})}{J1(\lambda_{1})^{2} + J0(\lambda_{1})^{2}} \cdot J0(\lambda_{1} \cdot 1) \cdot \exp(-\lambda_{1}^{2} \cdot \tau) = 0.578$$
$$\mathbf{T}_{\mathbf{S}} \coloneqq \mathbf{T}_{\text{inf}} + (\mathbf{T}_{0} - \mathbf{T}_{\text{inf}}) \cdot \theta_{\mathbf{S}} = 66.193 \cdot ^{\circ}\mathbf{C}$$